

## Addendum to 'The axial gauge for gravity'

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1981 J. Phys. A: Math. Gen. 14 3123

(<http://iopscience.iop.org/0305-4470/14/11/034>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

### Download details:

IP Address: 129.252.86.83

The article was downloaded on 31/05/2010 at 05:39

Please note that [terms and conditions apply](#).

**ADDENDUM**

**Addendum to ‘The axial gauge for gravity’**

R Delbourgo

Department of Physics, University of Tasmania, Box 252C, GPO, Hobart, Tasmania, Australia 7001

Received 7 September 1981

In a recent preprint Capper and Leibbrandt (1981) point to a discrepancy between their work and mine (Delbourgo 1981). While we agree that the graviton self-energy possesses  $n$ -dependent infinite parts they prove that it is not transversal, contrary to what Matsuki (1979) stated and what I purported to confirm in my Letter. The disagreement has prompted me to re-examine my calculations and search for a possible error. This I have now located and thus I now agree completely with Capper and Leibbrandt’s answers, even though I use the multiplier representation

$$\delta(n \cdot h) = \int db \exp(ibn \cdot h)$$

instead of their quadratic representation

$$\delta(n \cdot h) = \lim_{\alpha \rightarrow 0} \exp [i(n \cdot h)^2/\alpha].$$

The argument up to and including (7') in my Letter has no errors. However, whereas (7') does trivially lead to a transverse  $\Pi$  in Yang–Mills theory, the corresponding procedure *narrowly fails* for gravity because, in the multiplication of the third term of (7) by  $\Delta(q) \Delta(r)$ , there arise  $n$ -dependent pieces whose indices are not contracted against the propagator and thereby disappear (which they do in Yang–Mills). Instead, from identities like

$$\Delta_{\rho\sigma\mu\nu}^{-1}(q)\Delta^{\mu\nu\mu'\nu'}(q) = \frac{1}{2}(\delta_{\rho}^{\mu'}\delta_{\sigma}^{\nu'} + \delta_{\rho}^{\nu'}\delta_{\sigma}^{\mu'}) - \Delta_{\rho\sigma,\tau}^{-1}(q)\Delta^{\mu\nu,\tau}(q),$$

there survives the residual integral

$$p^{\kappa}\Pi_{\kappa\lambda,\kappa'\lambda'}(p) = -iK^2 \int \frac{p \cdot n}{q \cdot n} q^{\mu'} \Delta_{\lambda}^{\nu',\rho'\sigma'}(r) \Gamma_{\kappa'\lambda',\mu'\nu'\rho'\sigma'}(pqr) d^4q \neq 0.$$

This identically equals the ‘pincer graph’ contribution  $F$  in equation (3.11) of Capper and Leibbrandt, taken in their limit  $\alpha \rightarrow 0$ . *Note however that the absorptive part of  $\Pi$  stays transversal.*

The conclusion is thus even more dismal than I portrayed it in my Letter: although effectively ghost-free, axial gauge gravity is so complicated as to be practically useless; this is in sharp contrast with axial gauge chromodynamics.

**References**

- Delbourgo R 1981 *J. Phys. A: Math. Gen.* **14** L235  
Capper D M and Leibbrandt G 1981 *Phys. Lett.* **104B** 158  
Matsuki T 1979 *Phys. Rev. D* **19** 2879